The radiometer equation – a more careful look

In the lectures we derive the equation

$$SNR = \frac{T_A}{T_{sys}} \left(\Delta \nu \tau \right)^{1/2}, \qquad (1)$$

sometimes called the *radiometer equation*, to determine the signal-to-noise ratio from a source generating an antenna temperature T_A using a telescope with a system temperature T_{sys} , a bandwidth $\Delta \nu$ and an integration time τ . Although this result is rigorously correct we made a couple of hand-wavy approximations along the way, so here we'll do it more carefully.

A radio source (and any other contributing noise source) will generate a randomly fluctuating voltage u(t)at the antenna which we will assume has a gaussian probability distribution with a standard deviation σ and a mean of zero, so that the probability density for u is

$$p(u) = \frac{1}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-u^2}{2\sigma^2}\right).$$
 (2)

Our radio telescope system generates a signal V(t), proportional to $[u(t)]^2$, which we use as our estimator for the noise power. If the signal from the source is the only noise present, then the mean of *V* is proportional to the source's flux density. If there are other noise sources present then only part of *V* will be from the source. The pdf of *V* is one-sided ($V \ge 0$), and is related to p(u) by

$$p(V) |dV| = 2p(u) |du| \quad (V \ge 0).$$
 (3)

Using

$$\mathrm{d}V = 2u\,\mathrm{d}u,\tag{4}$$

we can write

$$p(V) = \frac{V^{-1/2}}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right).$$
 (5)

This is known as a *chisquared distribution with one degree of freedom*, and has a mean of

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$$\langle V \rangle = \int_0^\infty V p(V) \, \mathrm{d}V = \sigma^2,$$
 (6)

and a variance of

$$\operatorname{var}[V] = \langle V^2 \rangle - \langle V \rangle^2 \tag{7}$$

$$= \int_0^\infty V^2 p(V) \,\mathrm{d}V - \sigma^4 \tag{8}$$

$$= 2\sigma^4. \tag{9}$$

The standard deviation of *V* is therefore $\sqrt{2\sigma^2}$, or $\sqrt{2}\langle V \rangle$. If the random signal is the only 'noise' in the system, then the signal-to-noise ratio is simply

$$SNR = \frac{\langle V \rangle}{[\operatorname{var}(V)]^{1/2}} = \frac{1}{\sqrt{2}}.$$
 (10)

In the lectures we use a plausibility argument to say that this signal-to-noise ratio is ' \simeq 1', but the result above is more rigorous. If other noise sources are present things are a little more complicated, but the basic argument still holds. The signal will now be represented by the power from just the source, proportional to the antenna temperature T_A , and the total noise power by the system temperature T_{sys} , so that

$$SNR = \frac{T_A}{T_{sys}\sqrt{2}}.$$
 (11)

We can now consider how to improve this signal-to-noise ratio by integrating (averaging) the sample values of *V* over some interval of time τ . Let there be *N* samples of *V* (or *u*) in time τ . If the samples are statistically independent, and $N \gg 1$, then (by definition) the standard deviation of the average is \sqrt{N} less than that of a sample signal-to-noise ratio, i.e.,

$$SNR = \frac{T_A}{T_{svs}\sqrt{2}}\sqrt{N}.$$
 (12)

Samples u_i and u_j will be independent if

$$\langle u_i u_j \rangle = \langle u^2 \rangle \delta_{ij},\tag{13}$$

i.e., if the autocorrelation of u is a delta function at the origin. This corresponds (by the Wiener–Khinchin theorem) to a white ('flat') power spectrum for the noise. The Nyquist sampling theorem tells us that a sampling interval of τ/N corresponds to a bandwidth of Δv where

$$\Delta \nu = \frac{1}{2} \frac{N}{\tau},\tag{14}$$

SO

$$N = 2\Delta\nu\tau,\tag{15}$$

(in the lectures we say that ' $N \simeq \Delta \nu \tau$ '), and therefore our final signal-to-noise ratio is

SNR =
$$\frac{T_{\rm A}}{T_{\rm sys}2^{1/2}} (2\Delta\nu\tau)^{1/2} = \frac{T_{\rm A}}{T_{\rm sys}} (\Delta\nu\tau)^{1/2}$$
. (16)

Note that the two factors of $\sqrt{2}$ introduced by our more careful analysis cancel out, so the precise and 'approximate' results are the same.

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